





DOCTORAL COURSE XXXIII CYCLE - RESEARCH ACTIVITY REPORT I YEAR

# Correlation Plenoptic Imaging and Two-Level Quantum Systems

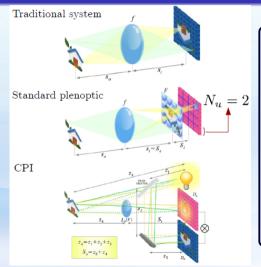
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Supervisors: Prof. Saverio Pascazio

#### **Outline**

- Correlation Plenoptic Imaging (CPI)
  - signal-to-noise properties of CPI with chaotic light
  - turbulence-free imaging
- Description of Jaynes-Cummings Model
  - Extension with Broken Inversion Symmetry
  - Different approaches to investigate

# Why Correlation Plenoptic Imaging CPI?



**Traditional optical imaging**: trade-off between resolution and depth of field

$$\Delta x = \frac{0.61\lambda}{NA}$$
  $DOF = \frac{\lambda}{NA^2}$ 

Standard plenoptic: only practical DOF gain

$$\Delta x = \frac{0.61\lambda}{NA}N_u$$
  $DOF = \frac{\lambda}{NA^2}N_u^2$ 

CPI: refocusing, 3D, no trade-off!

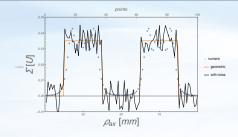
# Correlation Imaging, Signal-to-noise ratio

How to characterize in correlation imaging the noise due to the refocusing? **Signal-to-noise ratio** 

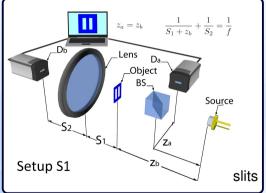
determines the optimal number of frames to obtain fast and high-quality images.

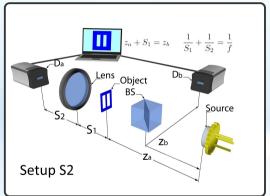
$$\sigma(\rho_a) = \int \Delta i_A (\alpha \rho_a + \beta \rho_b) \Delta i_B(\rho_b) d^2 \rho_b$$
Signal (image):  $\Sigma(\rho_a) = \langle \sigma(\rho_a) \rangle$ 
Variance:  $\mathcal{N} = \langle \sigma(\rho_a)^2 \rangle - \langle \sigma(\rho_a) \rangle^2$ 

 $\alpha, \beta$  refocusing parameters.



#### Signal-to-noise ration of CPI with chaotic light





1st order image of the source 2nd order image of the object.

1st order image of the object 2nd order image of the lens.

#### Trade-off SNR and resolution

The SNR at the point  $\rho_a$  of the detector  $D_a$ , whit  $N_f$  number of independent frames is

$$SNR(oldsymbol{
ho}_a) = \Sigma(oldsymbol{
ho}_a)\sqrt{rac{N_f}{\mathcal{N}}}$$

#### In preparation to be submitted to PRA

Setup 1: 
$$\frac{SNR(\rho_a)}{\sqrt{N_f}} \sim \sqrt{\frac{a^2}{A_{\text{obj}}}} \sqrt{\frac{\Delta x^2}{A_{\text{obj}}}} |A(\rho_a)|^2, \ \Delta x = (\lambda z/a)|1 - z_b/z_a|$$
Setup 2: 
$$\frac{SNR(\rho_a)}{\sqrt{N_f}} \sim \left(\frac{S_2/S_1}{1 - S_2/S_2^f}\right)^2 \sqrt{\frac{\sigma_B^2}{A_{\text{lens}}}} \frac{A_{\text{lens}}}{A_{\text{obj}}} \left|A\left(-\frac{\rho_a}{\mu}\right)\right|^2,$$

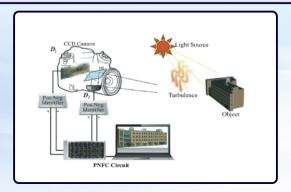
Setup 2: 
$$\frac{SNR(\rho_a)}{\sqrt{N_t}} \sim \left(\frac{S_2/S_1}{1-S_2/S_2^t}\right)^2 \sqrt{\frac{\sigma_B^2}{A_{lens}}} \frac{A_{lens}}{A_{obj}} \left| A\left(-\frac{\rho_a}{\mu}\right) \right|^2$$

a =linear size of the smallest transmissive part of the object.

### Turbulence-Free Imaging

Correlation Plenoptic Imaging

Second-order imaging is robust against refraction index fluctuations which usually introduce detrimental effects in the images.



Our aim is test to robustness against turbulence of CPI protocol.

Asymmetric Quantum Systems

# **Jaynes-Cummings Model**

The Jaynes-Cummings model (JCM) is a theoretical model in quantum optics for a system of a two-level atom interacting with a quantized mode of an optical cavity.

• 
$$\sigma^+ = |e\rangle\langle g|$$

• 
$$\sigma^- = |g\rangle\langle e|$$

 g<sub>r</sub> dipoletransition coupling constant.

$$H_{JC} = H_a + H_f + H_i$$

$$H_a = rac{\hbar \omega_a}{2} \sigma_z$$
 $H_f = \hbar \omega_c a^\dagger a$ 

$$H_I = g_r(\sigma^+ \otimes a + \sigma^- \otimes a^\dagger)$$

# **Dipole Approximation and Rotating Wave Approximation**

- for real atom a is of the order of Bohr radius
- the dominant wavelength is in resonance with the cavity  $(\omega_a \sim \omega_c)$

# **The Dipole Operator**

The dipole operator in our basis  $\{|e\rangle, |g\rangle\}$  is

$$\mathbf{d} = \begin{pmatrix} \mathbf{d}_{ee} & \mathbf{d}_{eg} \\ \mathbf{d}_{ge} & \mathbf{d}_{gg} \end{pmatrix} = (\mathbf{d}_{ee} + \mathbf{d}_{gg}) \frac{\mathbf{I}}{2} + (\mathbf{d}_{ee} - \mathbf{d}_{gg}) \sigma_z + \mathbf{d}_{eg} \sigma^+ + \mathbf{d}_{ge} \sigma^-.$$

The term  $d_{eg}$  is responsible for the transition between the energy levels.

What is the physical role of  $d_{ee}$ ?

#### Solvable Hamiltonian (without $d_{ee}$ )

The full Hamiltonian for symmetric system with no permanent dipole reads

$$H = \frac{\hbar \omega_{a}}{2} \sigma_{z} \otimes I_{\mathcal{F}} + \hbar \omega_{c} I_{\mathbb{C}^{2}} \otimes a^{\dagger} a + g_{r} (\sigma^{+} \otimes a + \sigma^{-} \otimes a^{\dagger})$$
$$= \begin{pmatrix} \hbar \omega_{c} (n-1) + \omega_{a}/2 & \sqrt{n} g_{r} \\ \sqrt{n} g_{r} & \hbar \omega_{c} n - \omega_{a}/2 \end{pmatrix}.$$

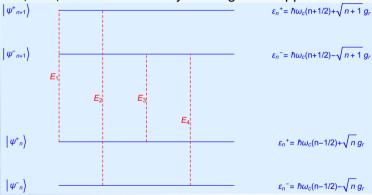
It commutes with the excitation number operator  $N = a^{\dagger} a + (\sigma_z + I)/2$  and its eigensystem is

$$|\psi_n^{\pm(0)}\rangle = C_{g,n}^{\pm}|g,n\rangle + C_{e,n}^{\pm}|e,n-1\rangle, \quad \varepsilon_n^{\pm(0)} = \hbar\omega_c\Big(n + \frac{1}{2}\Big) \pm \sqrt{\hbar^2 \frac{(\omega_c - \omega_a)^2}{4} + ng_r^2}$$

where  $C_{q,n}^{\pm}$ ,  $C_{e,n}^{\pm}$  are obtained by  $H|\psi_n^{\pm(0)}\rangle = \varepsilon_n^{\pm(0)}|\psi_n^{\pm(0)}\rangle$ .

#### **Possible Transition**

Transitions for  $n_f = n_i - 1$  are allowed by rotating wave approximation



$$E_{1,2} = \hbar \omega_c + g_r(\sqrt{n+1} \mp \sqrt{n}), \quad E_{3,4} = \hbar \omega_c - g_r(\sqrt{n+1} \mp \sqrt{n}).$$

#### Transitions at double frequencies of the cavity become opened

$$E_{1,2,3,4} = 2\hbar\omega_c \pm \hbar\frac{\Omega_R}{2}(\sqrt{n+1} \pm \sqrt{n})$$

$$E_{5,6,7} = \sqrt{n}\hbar\Omega_R$$

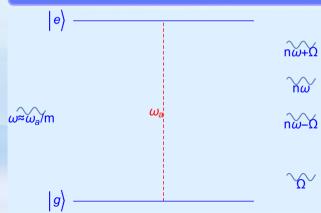
#### **Classical Interaction**

For the interaction with a classical electric field

$$H = rac{\hbar \omega_a}{2} \sigma_z \otimes \mathbf{I} - \mathbf{E} \cdot \mathbf{d} \cos(\omega_c t)$$

with wavefunction  $|\psi(t)\rangle = C_g(t)|g\rangle + C_e(t)|e\rangle$ , where  $C_g$ ,  $C_e$  are obtained by the Schrödinger equation, we can evaluate  $\mathbf{d}(t) = \langle \psi(t)|\mathbf{d}|\psi(t)\rangle$ .

$$\boldsymbol{d}(t) = \boldsymbol{d}_{ee} \frac{\Omega_R^2}{4\Omega^2} e^{i\Omega t} + \boldsymbol{d}_{eg} \frac{\Omega_R}{2\Omega} \sum_{n=-\infty}^{\infty} J_{m-n}(x_k) e^{in\omega t} \left[ \frac{\omega_a}{\Omega} + \frac{1}{2} \left( i - \frac{\omega_a}{\Omega} \right) e^{-i\Omega t} - \frac{1}{2} \left( i + \frac{\omega_a}{\Omega} \right) e^{i\Omega t} \right]$$



$$\bullet \Delta = \omega_c - m\omega$$



#### **Exams**

- Management and knowledge of European research model and promotion of research results by dr. D'Orazio;
- How to prepare a technical speech in English by prof. White;
- Programing with Python for Data Science by dr. Diacono;
- Introduction to C++ programming by dr. Cafagna;

- Linear Stability analysis by prof. Gonnella;
- Differential equations and physical phenomena by prof. Pascazio;
- States, observables and evolution by prof. Facchi;
- Atom-photon interactions by dr. Pepe.

#### References

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### acknowledgements

# Thank you for your attention any questions?

running over: "from turboless to turbulence-free" for comments and suggestions: giovanni.scala@ba.infn.it