



UNIVERSITÀ
DEGLI STUDI DI BARI
ALDO MORO



Waveguide QED

Domenico Pomarico

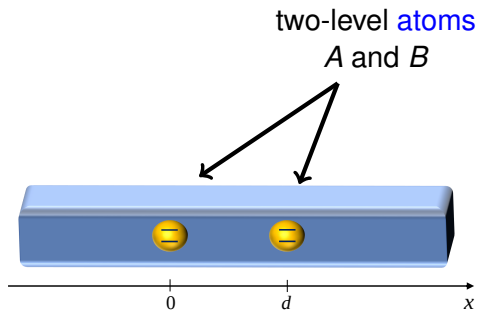
Corso di Dottorato XXXII ciclo - Relazione attività di ricerca II anno

16 Ottobre 2018

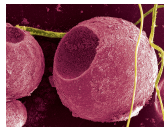
Supervisors: Prof. Saverio Pascazio, Dott. Francesco V. Pepe

Waveguide QED

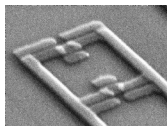
Multi-level quantum systems decay can be controlled by means of **boundary conditions** and **dimensional reduction** (e.g. embedding in a waveguide).



Quantum Dot



SQUID



Electromagnetic field quantization:

$$[a(k), a^\dagger(k')] = \delta(k - k')$$



just **photons** in the lowest-cutoff energy $TE_{1,0}$ mode

States & Dynamics

The Hamiltonian is parametrized by the **distance** d :

$$H_d = H_0 + \lambda V_d,$$

$$H_0 = H_S + H_B = \omega_0 (|e_A\rangle \langle e_A| + |e_B\rangle \langle e_B|) + \int dk \, \omega(k) a^\dagger(k) a(k),$$

and the dipolar interaction is expressed with **rotated couplings**:

$$V_d = \int dk \sum_{s=\pm} g_s(k) b_s a^\dagger(k) + h.c. \quad \text{with} \quad g_s(k) = \frac{g_0(k) + s g_d(k)}{\sqrt{2}},$$

where $b_s^\dagger |g_A, g_B\rangle = \frac{b_A^\dagger + s b_B^\dagger}{\sqrt{2}} |g_A, g_B\rangle = |\Psi^s\rangle$, $b^\dagger = |e\rangle \langle g|$, $g_x(k) = \frac{e^{ikx}}{\sqrt{\omega(k)}}$.

States & Dynamics

The Hamiltonian is parametrized by the **distance** d :

$$H_d = H_0 + \lambda V_d,$$

$$H_0 = H_S + H_B = \omega_0 (|e_A\rangle \langle e_A| + |e_B\rangle \langle e_B|) + \int dk \omega(k) a^\dagger(k) a(k),$$

and the dipolar interaction is expressed with **rotated couplings**:

$$V_d = \int dk \sum_{s=\pm} g_s(k) b_s a^\dagger(k) + h.c. \quad \text{with} \quad g_s(k) = \frac{g_0(k) + s g_d(k)}{\sqrt{2}},$$

where $b_s^\dagger |g_A, g_B\rangle = \frac{b_A^\dagger + s b_B^\dagger}{\sqrt{2}} |g_A, g_B\rangle = |\Psi^s\rangle$, $b^\dagger = |e\rangle \langle g|$, $g_x(k) = \frac{e^{ikx}}{\sqrt{\omega(k)}}$.

Bell basis

States & Dynamics

The Hamiltonian is parametrized by the **distance** d :

$$H_d = H_0 + \lambda V_d,$$

$$H_0 = H_S + H_B = \omega_0 (|e_A\rangle \langle e_A| + |e_B\rangle \langle e_B|) + \int dk \omega(k) a^\dagger(k) a(k),$$

and the dipolar interaction is expressed with **rotated couplings**:

$$V_d = \int dk \sum_{s=\pm} g_s(k) b_s a^\dagger(k) + h.c. \quad \text{with} \quad g_s(k) = \frac{g_0(k) + s g_d(k)}{\sqrt{2}},$$

where $b_s^\dagger |g_A, g_B\rangle = \frac{b_A^\dagger + s b_B^\dagger}{\sqrt{2}} |g_A, g_B\rangle = |\Psi^s\rangle$, $b^\dagger = |e\rangle \langle g|$, $g_x(k) = \frac{e^{ikx}}{\sqrt{\omega(k)}}$. **Bell basis**

The **number of excitations** $\mathcal{N}$ is a constant of motion, so for $\mathcal{N} = 2$:

$$|\psi\rangle = c_{AB} |e_A, e_B\rangle + \sum_{s=\pm} \int dk B_s(k) |\Psi^s; k\rangle + \frac{1}{\sqrt{2}} \int dk_1 dk_2 A(k_1, k_2) |k_1, k_2\rangle$$

with $B_s \in L^2(\mathbb{R})$, $A \in L^2(\mathbb{R})^{\otimes 2}$ symmetric for $k_1 \leftrightarrow k_2$.

Facchi, Pascazio, Pepe, Pomarico, in submission PRA, AX11558

Perturbative expansion

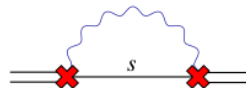
$$\omega(k) = \omega(-k)$$

$$g_s(k) = g_s^*(-k)$$

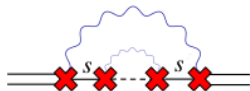


Parity conservation for single excitation propagators linked by a photon

Double excitation propagator
lower order contributions



→ $O(\lambda^2)$ useful to estimate **plasmon** eigenstate energy



→ planar $\left\{ \int dq g_{s'}^*(q) g_s(q) F_{ss'}(z, \omega(q)) \right.$



→ crossed $\left\{ g_{\pm}^*(q) g_{\mp}(q) \propto \sin qd \implies s = s' \right.$



→ no parity constraint

Renormalized double excitation propagator

$$\begin{aligned}
 \langle e_A, e_B | \mathcal{G}(z) | e_A, e_B \rangle &= \text{diagram} = \frac{1}{z - 2\omega_0 - \lambda^2 \Sigma_2(z)} = \mathcal{G}_2(z)
 \end{aligned}$$

The diagram shows a double line with a blue bar on the left, equal to a double line with a blue bar on the right, plus a diagram with a double line, a red cross, a blue circle labeled $\Sigma_2(z)$, another red cross, and a blue bar on the right.

Renormalized single excitation propagator (one-excitation sector):

$$\mathcal{G}_s(z) = \text{diagram} = \text{diagram} + \text{diagram}$$

The diagram shows a blue circle on a line, equal to a line, plus a diagram with a line, a red cross, a dashed line labeled $\Sigma_s(z)$, another red cross, and a blue circle on a line.

the vertex renormalization is obtained:

$$\lambda s g_s(k)(1 + X_s(z, k)) = \text{diagram} = \text{diagram} + \text{diagram}$$

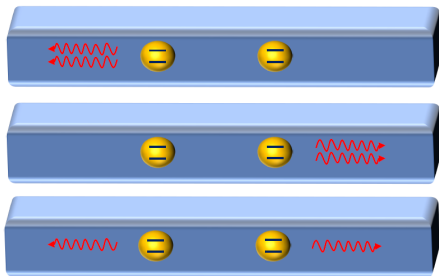
The diagram shows a red cross on a blue bar, equal to a red cross, plus a diagram with a red cross, a dashed line, a blue circle, and a red cross on a blue bar. A wavy line labeled k connects the two red crosses in the second diagram.

Two-photon amplitude

Asymptotic **two photon** state:

$$\langle k_1, k_2 | e^{-iHt} | e_A, e_B \rangle = A(k_1, k_2, t) \xrightarrow{t \rightarrow \infty} \mathcal{A}(k_1, k_2) e^{-i(\omega(k_1) + \omega(k_2))t}$$

$$D_{\Leftrightarrow [=]} = \{(k_1, k_2) \in \mathbb{R}^2 | \text{sgn}(k_1) = [-] \text{sgn}(k_2)\}$$



$$P_{\Leftrightarrow} = \frac{1}{2} \int_{D_{\Leftrightarrow}} dk_1 dk_2 |\mathcal{A}(k_1, k_2)|^2$$

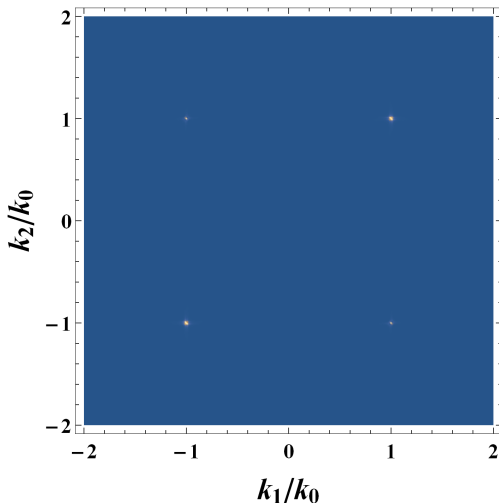
$$P_{\Rightarrow} = \frac{1}{2} \int_{D_{\Rightarrow}} dk_1 dk_2 |\mathcal{A}(k_1, k_2)|^2$$

Ratio of asymptotic
probabilities

$$R(\lambda) \Big|_{\lambda \rightarrow 0} = \frac{P_{\Leftrightarrow}}{P_{\Rightarrow}} = 1 + 2 \sin^2 k_0 d$$

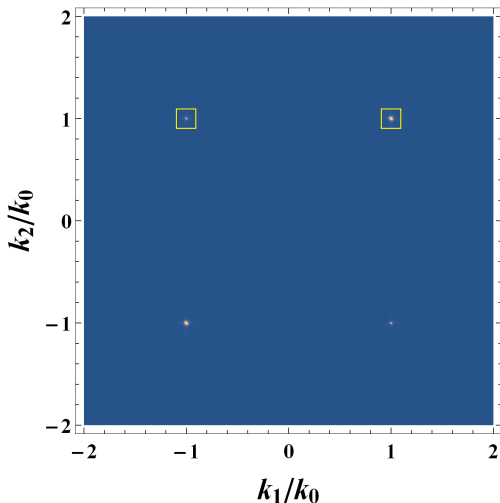
Non-perturbative weak coupling limit

Selected parameters: $\omega_0 = 1.1M$, $\lambda = 10^{-2}M$, $k_0 d = \frac{\pi}{2}$



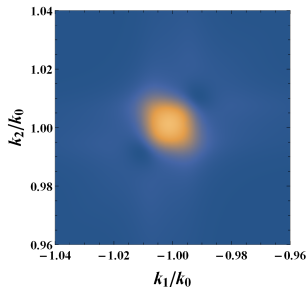
Non-perturbative weak coupling limit

Selected parameters: $\omega_0 = 1.1M$, $\lambda = 10^{-2}M$, $k_0 d = \frac{\pi}{2}$

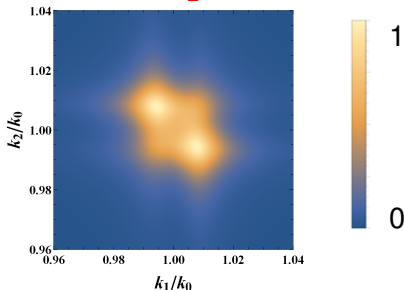


Non-perturbative weak coupling limit

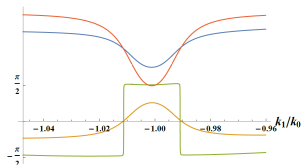
Selected parameters: $\omega_0 = 1.1M$, $\lambda = 10^{-2}M$, $k_0 d = \frac{\pi}{2}$



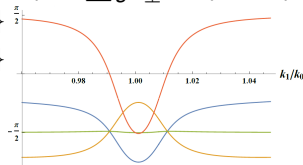
$$\frac{k_0}{M} \simeq 0.458$$

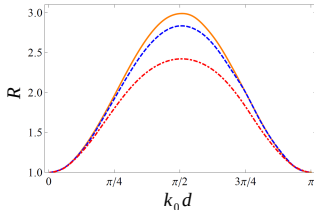


This is the **parity contributions interference** $\mathcal{A}(k_1, k_2) = \sum_{s=\pm} \mathcal{A}_s(k_1, k_2)$:

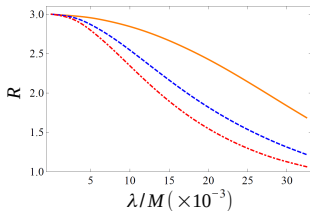


- $\phi_- = \text{Arg}\{\mathcal{A}_-(k_1, k_2)\}$
- $\phi_+ = \text{Arg}\{\mathcal{A}_+(k_1, k_2)\}$
- $\text{Arg}\{\mathcal{A}(k_1, k_2)\}$
- $\phi_+ - \phi_-$



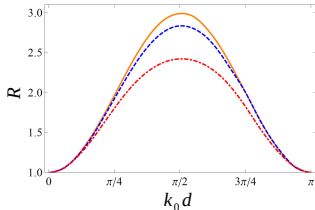


- $\lambda = 10^{-3}M$
- - $\lambda = 10^{-2}M$
- . $\lambda = 2 \times 10^{-2}M$

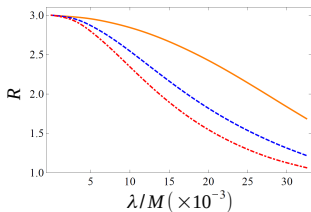


- $k_0 d = \pi/2$
- - $k_0 d = 3\pi/2$
- . $k_0 d = 5\pi/2$

- $\omega_0 = 1.1M$;
- peaks in the variation of $k_0 d$ are points of the **orange** curve in the variation of λ ;
- increased distances implement stronger coupling regimes.



- $\lambda = 10^{-3}M$
- - $\lambda = 10^{-2}M$
- ... $\lambda = 2 \times 10^{-2}M$

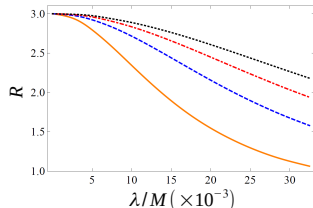


- $k_0 d = \pi/2$
- - $k_0 d = 3\pi/2$
- ... $k_0 d = 5\pi/2$

- $k_0 d = 5\pi/2$;
- increasing values for ω_0 push forward the plasmonic bound state.

- $\omega_0 = 1.1M$ —
- $\omega_0 = 1.2M$ - -
- $\omega_0 = 1.3M$...
- $\omega_0 = 1.4M$..

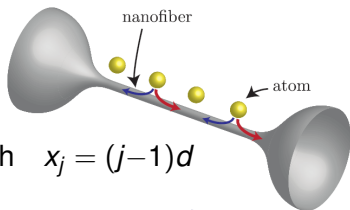
- $\omega_0 = 1.1M$;
- peaks in the variation of $k_0 d$ are points of the orange curve in the variation of λ ;
- increased distances implement stronger coupling regimes.



Multi-atom chain

$$H_0 = \sum_{j=1}^N \omega_0 b_j^\dagger b_j + \int dk \omega(k) a^\dagger(k) a(k),$$

$$V_d = \sum_{j=1}^N \int dk \left(g_{x_j}(k) b_j^\dagger a(k) + h.c. \right), \quad \text{with} \quad x_j = (j-1)d$$



Excitation amplitude array $\mathbf{c} = (c_1, \dots, c_N)$, self-energy matrix $\Sigma(z)$

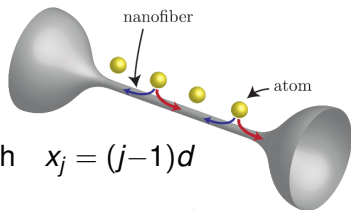
Hamiltonian eigenvalues ($E \in \mathbb{R}$): $((\omega_0 - E)\mathbb{1} - \Sigma(E + i0^+)) \cdot \mathbf{c} = 0$

Waveguide QED: $\Sigma(z) = \Sigma_{\text{pole}}(z) + \Sigma_{\text{cut}}(z)$

Multi-atom chain

$$H_0 = \sum_{j=1}^N \omega_0 b_j^\dagger b_j + \int dk \, \omega(k) a^\dagger(k) a(k),$$

$$V_d = \sum_{j=1}^N \int dk \, \left(g_{x_j}(k) b_j^\dagger a(k) + h.c. \right), \quad \text{with} \quad x_j = (j-1)d$$



Excitation amplitude array $\mathbf{c} = (c_1, \dots, c_N)$, self-energy matrix $\Sigma(z)$

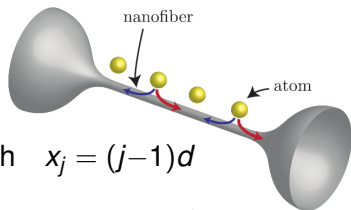
Hamiltonian eigenvalues ($E \in \mathbb{R}$): $((\omega_0 - E)\mathbb{1} - \Sigma(E + i0^+)) \cdot \mathbf{c} = 0$

Waveguide QED: $\Sigma(z) = \Sigma_{\text{pole}}(z) + \cancel{\Sigma_{\text{cut}}(z)} \longleftarrow \mathcal{O}(e^{-Md})$

Multi-atom chain

$$H_0 = \sum_{j=1}^N \omega_0 b_j^\dagger b_j + \int dk \omega(k) a^\dagger(k) a(k),$$

$$V_d = \sum_{j=1}^N \int dk \left(g_{x_j}(k) b_j^\dagger a(k) + h.c. \right), \quad \text{with } x_j = (j-1)d$$



Excitation amplitude array $\mathbf{c} = (c_1, \dots, c_N)$, self-energy matrix $\Sigma(z)$

Hamiltonian eigenvalues ($E \in \mathbb{R}$): $((\omega_0 - E)\mathbb{1} - \Sigma(E + i0^+)) \cdot \mathbf{c} = 0$

Waveguide QED: $\Sigma(z) = \Sigma_{\text{pole}}(z) + \cancel{\Sigma_{\text{cut}}(z)} \leftarrow \mathcal{O}(e^{-Md})$

Resonance condition $\sqrt{E^2 - M^2}d = n\pi$

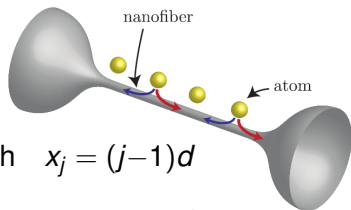
n even: $\sum_j c_j = 0$, n odd: $\sum_j (-1)^j c_j = 0$

Facchi, Lonigro, Pascazio, Pepe, Pomarico, in preparation

Multi-atom chain

$$H_0 = \sum_{j=1}^N \omega_0 b_j^\dagger b_j + \int dk \omega(k) a^\dagger(k) a(k),$$

$$V_d = \sum_{j=1}^N \int dk \left(g_{x_j}(k) b_j^\dagger a(k) + h.c. \right), \quad \text{with } x_j = (j-1)d$$



Excitation amplitude array $\mathbf{c} = (c_1, \dots, c_N)$, self-energy matrix $\Sigma(z)$

Hamiltonian eigenvalues ($E \in \mathbb{R}$): $((\omega_0 - E)\mathbb{1} - \Sigma(E + i0^+)) \cdot \mathbf{c} = 0$

Waveguide QED: $\Sigma(z) = \Sigma_{\text{pole}}(z) + \Sigma_{\text{cut}}(z) \rightarrow$ parity metastability with first decaying state according to:

Resonance condition $\sqrt{E^2 - M^2}d = n\pi$

n even: $\sum_j c_j = 0$, n odd: $\sum_j (-1)^j c_j = 0$

$$\frac{\gamma}{2} \propto \lambda^2 \frac{2(N-2)}{3N^3}$$

Facchi, Lonigro, Pascazio, Pepe, Pomarico, in preparation

Outlook

Main activities in progress for the third year are:

- variation of the **initial state** with two excitations;
- existence of **non resonant eigenstates** for the multi-atom setup;
- investigation of waveguides endowed with **discrete spectrum**;
- dissipation control in density matrix **master equation**;
- application for the simulation of the **Heisenberg model**.

During the second year I attended the following conferences and summer school:

- "Statistical Mechanics & Field Theory", Bari;
- "Problemi Attuali di Fisica Teorica", Vietri sul mare, with talk "Bound states and entanglement generation in waveguide QED";
- "Information Geometry, Quantum Mechanics and Applications", San Rufo-Policeta, with talk "Correlated photon emission by two excited atoms in a waveguide";
- "Collective behavior in quantum matter", summer school in ICTP (Trieste), with poster "Correlated photon emission by two excited atoms in a waveguide".